

1- Recall the Lagrangian you used to derive geodesic equations. Show that  $\mathcal{L}$  is constant along any geodesics. i.e. that

$$\frac{d}{d\tau} \left( g_{\mu\nu} \frac{dx^\mu}{d\tau} \frac{dx^\nu}{d\tau} \right) = 0$$

for  $x^\mu(\tau)$  a solution of the geodesic equation. What is the symmetry that by Noether's theorem, gives rise to this constant of motion?

2 -The covariant derivatives of a covector field  $A_\mu$  and a  $(0, 2)$  tensor  $B_{\mu\nu}$  field are:

$$\begin{aligned} \nabla_\mu A_\nu &= \partial_\mu A_\nu - \Gamma_{\mu\nu}^\lambda A_\lambda \\ \nabla_\lambda B_{\mu\nu} &= \partial_\lambda B_{\mu\nu} - \Gamma_{\lambda\mu}^\rho B_{\rho\nu} - \Gamma_{\lambda\nu}^\rho B_{\mu\rho} \end{aligned} \quad (1)$$

(a) Show that, even though  $\partial_\mu A_\nu$  is not a tensor (which is why the  $\Gamma$ -term is required in  $\nabla_\mu A_\nu$ ), the curl (or rotation)  $\partial_\mu A_\nu - \partial_\nu A_\mu$  is a tensor, then show that the covariant curl of a covector is equal to its ordinary curl,

$$\nabla_\mu A_\nu - \nabla_\nu A_\mu = \partial_\mu A_\nu - \partial_\nu A_\mu$$

This provides an alternative argument for the fact that  $\partial_\mu A_\nu - \partial_\nu A_\mu$  is a tensor.

(b) Show that the covariant derivative of the metric is zero,  $\nabla_\lambda g_{\mu\nu} = 0$ .

3- Establish

$$\Gamma_{\alpha\beta}^{\prime\mu} = \frac{\partial x^{\prime\mu}}{\partial x^\nu} \frac{\partial x^\sigma}{\partial x^{\prime\alpha}} \frac{\partial x^\rho}{\partial x^{\prime\beta}} \Gamma_{\sigma\rho}^\nu - \frac{\partial x^\nu}{\partial x^{\prime\alpha}} \frac{\partial x^\sigma}{\partial x^{\prime\beta}} \frac{\partial^2 x^{\prime\mu}}{\partial x^\nu \partial x^\sigma} \quad (2)$$

by assuming that the quantity defined by

$$\nabla_\mu A^\alpha = \partial_\mu A^\alpha + \Gamma_{\nu\mu}^\alpha A^\nu \quad (3)$$

has the tensor character indicated. Take the partial derivative of

$$\delta_\mu^{\prime\alpha} = \frac{\partial x^{\prime\alpha}}{\partial x^\mu} = \frac{\partial x^{\prime\alpha}}{\partial x^\beta} \frac{\partial x^\beta}{\partial x^\mu}$$

with respect to  $x^{\prime\nu}$  to establish the alternative form:

$$\Gamma_{\alpha\beta}^{\prime\mu} = \frac{\partial x^{\prime\mu}}{\partial x^\nu} \frac{\partial x^\sigma}{\partial x^{\prime\alpha}} \frac{\partial x^\rho}{\partial x^{\prime\beta}} \Gamma_{\sigma\rho}^\nu + \frac{\partial x^{\prime\mu}}{\partial x^\nu} \frac{\partial^2 x^\nu}{\partial x^{\prime\alpha} \partial x^{\prime\beta}} \quad (4)$$

2 and 4 are equivalent.

4- [d'inverno 6.5] assuming 3 and  $\nabla_\alpha \phi = \partial_\alpha \phi$  apply the Leibniz rule to covariant derivative of  $A_\mu A^\mu$ , where  $A^\mu$  is arbitrary to verify 1.

5- [Carroll 3.5] Consider a 2-Sphere with coordinates  $(\theta, \phi)$  and metric:

$$ds^2 = d\theta^2 + \sin^2 \theta d\phi^2$$

Show that the line of constant longitude ( $\phi = \text{constant}$ ) are geodesics, and that the only line of constant latitude ( $\theta = \text{constant}$ ) that is a geodesic is the equator ( $\theta = \pi/2$ ).